

Propagation properties of dissipative waves in non Maxwellian plasmas

Abstract

The propagation of the ion-acoustic shock waves (dissipative waves) are presented in a plasma consisting of fluid ions and cold and hot electrons. Cold and hot electrons are considered as Maxwellian and kappa distribution functions, respectively. Using reductive perturbation technique, Korteweg-de Vries-Burger (KdVB) and Burgers equations are derived for shock waves in the present plasma model. It is found that both positive (compressive) and negative (rarefactive) ion-acoustic shock waves can be propagating in present medium. Finally, the influence of different parameters on both these shock waves are studied numerically. Our results are important for the electrostatic shock wave structures in space plasmas.

Keywords: two-temperature plasmas, K-dVB equation, burgers equation, ion-acoustic shock waves, superthermal electrons

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Introduction

One of the most important subjects in the plasma physics is the propagation of nonlinear waves that can be described by different partial differential equations as Korteweg-de Vries (K-dV) equation,¹ Korteweg-de Vries Burger (K-dVB) equation² and the nonlinear Schrödinger equation (NLSE).³ Electrostatic solitary waves (solitons) that described usually by K-dV equation derived by balancing the nonlinearity of the medium with the wave dispersion. On the other hand, for a dissipative media, the balance between nonlinearity and dissipation may lead to formation of the shock structure. Shock waves in plasmas are different from these in normal gases. Because of presence of charge particles, certain types of shock waves can exist in a plasma. Shock waves can also be observed in space⁴⁻⁶ and experimental⁷⁻⁹ plasmas. Plasma with two-distinct group of electrons (cold and hot) can also be seen in space plasmas.^{10,11} Many papers have investigated the propagation of ion-acoustic solitons in a plasma with two types electrons.¹²⁻¹⁴ In such as media the theoretical and experimental results have been shown that the characteristics of solitary waves strongly modified by the presence of minority population of cold electrons.¹⁵⁻¹⁷

However, in the space and astrophysical plasma environments such as solar wind, magnetosphere and auroral zone plasmas, particles are often characterized by velocity distributions that are not Maxwellian and modeled by a kappa (or generalized Lorentzian) distribution function.¹⁸⁻²⁰ Recently, the propagation of ion acoustic solitons in collisionless plasmas with high-energy particles (represented by kappa distribution) is investigated in many papers.²¹⁻²⁷ In addition, the existence and propagation of ion acoustic solitons in a plasma with two-temperature kappa distributed electrons have been investigated in.²⁸⁻³⁰ They have been studied treatment of small and arbitrary solitary waves by deriving an energy integral equation. On the other hand, it should be remarked that a medium with dispersive and significant dissipative properties supports another solitary wave such as shock waves instead of solitons. Many researchers have studied shock wave structures in superthermal plasmas. For example, Sultana et al.³¹ have studied dynamic of shock waves in a collisionless electron-ion plasma with superthermal electrons which is modeled

by kappa distribution function. They obtained that higher deviations from a pure Maxwellian behavior induced shocks which were larger, narrower and faster.³¹ Ion-acoustic shock waves in a plasma consisting of electrons, positrons and ions with Boltzmann distributed positrons and superthermal electrons were also studied by Pakzad.³²

It is observed that an increasing positron concentration decreases the amplitude of the waves. The characteristics of ion-acoustic shock waves in electron-positron-ion plasma were investigated in³³ where electrons and positrons were modeled by kappa distribution function. It is seen that the amplitude of ion-acoustic shock waves decreases as the positron density and electron temperature increases.³³ In addition, Shah et al.³⁴ have investigated the effect of positron beam on the ion-acoustic shock wave in a plasma consisting of inertial ions and superthermal electrons and the positron beam. They have derived the spectral index of the superthermal electrons, and concentration of impinging positron beam have significant effects on both the amplitude and steepness of the ion-acoustic shock wave. The oblique propagation of ion-acoustic solitons in magnetized plasmas containing kappa distributed hot and cold electrons have examined by several authors.^{35,36} In addition, Bains et al.³⁷ studied the propagation of ion-acoustic shock waves in a magnetized plasma consisting of cold ions and two different temperature kappa distributed electrons. It is seen that all plasma parameters such as the density of hot electrons, the superthermality of electrons, the obliqueness, the strength of external magnetic field and the ion kinematic viscosity increase the shock height. The investigation of dust ion acoustic shock waves in two temperature plasmas has attracted significant attention among researchers.³⁸⁻⁴⁰ However, the properties of shock waves in a plasma containing cold ions, cold (Maxwellian) and hot (superthermal) electrons has not been reported so far in papers. Therefore, in this paper using reductive perturbation technique, we have investigated shock structures from the solutions of the K-dV Burgers equation. Both the cases with quasi-neutrality condition as well as Poisson's equation have been considered to observe such type of structures. It has been noticed that with the presence of the dissipation terms, a significant variation of the structures is observed. The layout of this paper is as follows: in section 2, we present basic equations and obtain K-dVB and Burgers equations. Also, these equations are solved by

tanh method. Numerical results are given in Sec. 3. Finally, in Sec. 4 the conclusions of the present work are discussed.

Basic equation

We consider one nonlinear three -component homogeneous plasma consisting of cold fluid ions, cold and hot (superthermal) electrons. Here, it is assumed that cold and hot electrons have Maxwellian and kappa distribution functions, respectively. In equilibrium, the quasi-neutrality condition is as $n_{c0} + n_{h0} = n_{i0}$, where n_{c0} , n_{h0} and n_{i0} are the unperturbed number densities of the cold electrons, hot electrons and ions, respectively. It is assumed that the present medium has dissipative properties in addition to what mentioned earlier. The fluid model for the ions is described by the following continuity and momentum equations:

$$\frac{\partial n_i}{\partial t} + \frac{\partial}{\partial x}(n_i u_i) = 0, \tag{1}$$

$$m_i n_i \frac{\partial u_i}{\partial t} - m_i n_i u_i \frac{\partial u_i}{\partial x} + e n_i \frac{\partial \Phi}{\partial x} - m_i n_i \frac{\partial^2 u_i}{\partial x^2} = 0. \tag{2}$$

The system is closed by Poisson’s equation:

$$\frac{\partial^2 \Phi}{\partial x^2} = 4\pi e(n_c + n_h - n_i). \tag{3}$$

In the above equation n_c , n_h and n_i are the density of cold electrons, hot electrons and ions, respectively. u_i is the velocity of ions; and Φ is the electrostatic potential. An *ad hoc* damping term is introduced in the momentum equation, involving the (ion) kinematic viscosity η_i . e is the magnitude of the electron charge and m_i is the ion mass. In our analysis, all parameters are normalized as follows, the ion and cold (hot) electron densities by the unperturbed ion density n_{i0} , the ion velocity by $c_s = (k_B T_e / m_i)^{1/2}$, the electrostatic potential by $\phi_0 = (k_B T_e / e)$, the space and time are normalized by Debye length $\lambda_D = C_s / \omega_{pi}$ and ion plasma frequency $\omega_{pi} = (e^2 n_{i0} / \epsilon_0 m_i)^{1/2}$ respectively. Therefore, the basic normalized equations in this system can be written as:

$$\frac{\partial n}{\partial t} + \frac{\partial}{\partial x}(nu) = 0, \tag{4}$$

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + \frac{\partial \phi}{\partial x} - \eta \frac{\partial^2 u}{\partial x^2} = 0, \tag{5}$$

$$\frac{\partial^2 \phi}{\partial x^2} - n_c + n_h - n = 0. \tag{6}$$

Where the normalized (ion) kinematic viscosity variable is defined as $\eta = \eta_i / (\omega_{pi} \lambda_D^2)$. The normalized number densities of electrons are given by

$$n_c = \alpha e^{\sigma \phi} \quad n_h = \beta \left(1 - \frac{S \phi}{k - 3/2} \right)^{-k+1/2}, \tag{7}$$

Where $\alpha = n_{c0} / n_{i0}$, $\beta = n_{h0} / n_{i0}$, $S = T_e / T_h = [(1 - \beta)\tau + \beta]^{-1}$ and $\sigma = T_e / T_c = \tau S$, where T_c and T_h are the temperature of cold and hot electrons, respectively and the variable τ is defined as hot to cold electron temperature ratio, namely $\tau = T_h / T_c$. The effective temperature T_e given as

$$T_e = T_c T_h \left[\left(\frac{n_{c0}}{n_{i0}} \right) T_h + \left(\frac{n_{h0}}{n_{i0}} \right) T_c \right]^{-1}. \tag{8}$$

Here k is the spectral index that measuring deviation from Maxwellian equilibrium and valid for $k > 3/2$.

Ion-Acoustic shock waves

We are now interested for investigation propagation of shock waves in present plasma model. Therefore, we shall employ reductive perturbation technique.⁴¹ The independent variables can be stretched as

$$\xi = \epsilon^{1/2}(x - \lambda t) \quad \text{and} \quad \tau = \epsilon^{3/2}t, \tag{9}$$

Where $\epsilon \ll 1$ is a small parameter and λ is the unknown linear phase velocity to be determined later. Also, the independent variables, and can expanded as follow

$$\begin{pmatrix} n \\ u \\ \phi \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} + \sum_{l=0}^{\infty} \epsilon^{l+1} \begin{pmatrix} n_{l+1} \\ u_{l+1} \\ \phi_{l+1} \end{pmatrix}. \tag{10}$$

In continue, substituting equations (9) and (10) into equations (4)–(6), and isolating distinct orders in ϵ , we obtain a Korteweg-de Vries Burgers (K-dVB) type equation in the form

$$\frac{\partial}{\partial \tau} \phi_1 + A \phi_1 \frac{\partial}{\partial \xi} \phi_1 + B \phi_1 \frac{\partial^3 \phi_1}{\partial \xi^3} - C \frac{\partial^3 \phi_1}{\partial \xi^2} = 0 \tag{11}$$

Where

$$\begin{aligned} A &= - \left[\beta C_2 S^2 + \frac{1}{2}(1 - \beta)\sigma^2 \right] \lambda^3 + \frac{3}{2\lambda}, \\ B &= \frac{\lambda^3}{2}, \\ C &= \frac{\eta_0}{2}, \end{aligned} \tag{12}$$

Where

$$\begin{aligned} \lambda &= \sqrt{\frac{1}{\alpha \sigma + \beta C_1 S}}, \\ C_1 &= \frac{\left(k - \frac{1}{2} \right)}{\left(k - \frac{3}{2} \right)}, \quad C_2 = \frac{\left(k - \frac{1}{2} \right) \left(k + \frac{1}{2} \right)}{2 \left(k - \frac{3}{2} \right)^2}. \end{aligned} \tag{13}$$

Here $\eta = \epsilon^{1/2} \eta_0$ where η_0 is the finite quantity of the order of unitary. On the other hand, if we assume the quasi-neutrality condition by employing the reductive perturbation method and using equations (4)–(5) and the Poisson’s equation in form

$$0 = \alpha e^{\sigma \phi} - \beta \left(1 - \frac{S \phi}{k - 3/2} \right)^{-k+1/2} - n, \tag{14}$$

one can drive the Burgers equation as follow:

$$\frac{\partial}{\partial \tau} \phi_1 + A \phi_1 \frac{\partial}{\partial \xi} \phi_1 + C \frac{\partial^2 \phi_1}{\partial \xi^2} = 0, \tag{15}$$

Finally, we will obtain exact solution for ion acoustic shock waves in current plasma model by tanh method.^{42,43} For K-dVB equation (11) a monotonic, kink-like shaped, solution is obtained via using this method in the form:

$$\phi_1 = \frac{3C^2}{25AB} \left[4 - (1 + \tanh X)^2 \right], \tag{16}$$

Where we have performed a transformation of the coordinates $X = k(\xi - V\tau)$; here $k = C/(10B)$ and $V = 6C^2/(25B)$ are related to the inverse width and the speed of the shock, respectively.^{42,43}

In which for Burger equation (15) we have

$$\phi_1 = \phi m \left[1 - \tanh(\zeta / \Delta) \right], \tag{17}$$

where the amplitude ϕm and the width Δ are given by

$$\phi m = \frac{U}{A}, \text{ and } \Delta = \frac{2C}{U}. \tag{18}$$

(Here $\zeta = k(\xi - U\tau)$, and $U = C^2/(5B)$ is a velocity [42-43].

Numerical results

In this study, a rigorous theoretical investigation of electrostatic shock structures in two electron temperature plasma (having cold and superthermal electron) with cold ions are carried out. By employing

reductive perturbation method, K-dV Burgers equation is derived. It is found that the coefficients of the K-dVB equation are significantly modified by ratio of unperturbed density of hot electron to ion ($\beta = n_{h0} / n_{i0}$), hot to cold electron temperature ratio ($\tau = T_h / T_c$) and superthermality parameter k . By assuming quasineutrality, Burgers equation is also derived and its details will be investigated in the following.

It should be noted that the nonlinear coefficient A have negative, zero and positive values. Therefore, it is observed that shock potential profiles are positive (negative) if $A > 0$ ($A < 0$). On the other hand, the vanishing of the nonlinear coefficient in Eqs. (11) and (15) (i.e., when $A = 0$) determines the critical values of $\beta = \beta_c$. In continue, to find the parametric regimes for the existence of opposite polarity shock potential profiles, we have numerically analyzed the nonlinear term A and obtain $A = 0$ (β versus τ) curve in Figure 1. From this figure, it is clear that for the fixed values of $\tau = T_h / T_c = 20$ and $k = 3$, we have two critical values 0.812 and 0.976, where for densities lie in two ranges $0 < \beta < \beta_{c1}$ ($= 0.812$) and $0 < \beta < \beta_{c2}$ ($= 0.976$), the nonlinear coefficient is positive and we have compressive shock waves, while the densities lie between these two critical values, namely $\beta_{c1} < \beta < \beta_{c2}$, lead to the rarefactive shock waves.

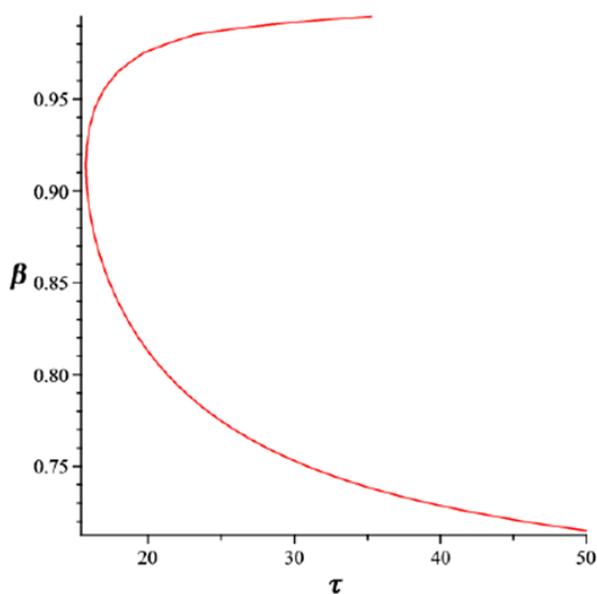


Figure 1 The showing variations of the nonlinear coefficient $A = 0$, the number density β versus temperature τ for $k = 3$.

Potential profiles of the ion acoustic shock waves, as given by Eq. (16), for fixed values $\tau = 20, k = 3$ and different cases of density β are investigated in Figure 2. It is observed that the amplitude of positive shock waves increases (decreases) with β for case $0 < \beta < \beta_{c1}$ ($\beta < \beta_{c2}$), (see Figure 2(a) and 2(c), respectively) in which for $\beta_{c1} < \beta < \beta_{c2}$ (rarefactive shock waves) the amplitude decreases by increasing the hot electron density β see Figure 2(b). It should be remarked that our results are different from results given in.³⁶ They studied the influence of hot electron density population on shock waves and shown that there is only a critical value ≈ 0.58 and in the both cases (compressive and rarefactive shock waves) the amplitude of shock increases with increase in the hot electron density.³⁷

The hot to cold electron temperature ratio τ is examined in Figure 3. It is shown that for a fixed value of density $\beta = 0.7$ (where it is

assumed lower than first critical density β_{c1}) the amplitude of the shock waves increases as τ increases (see Figure 3(a)), while from Figure 3(b) it is found that for case $\beta = 0.85$ (between two critical value) by increasing hot electron temperature, the amplitude of shock waves decreases. Finally, from Figure 3(c), it is clear that for case $\beta = 0.985$ (bigger than second critical density β_{c2}) the amplitude increases as τ increases. This means with increase temperature the plasma model can support larger (smaller) compressive (rarefactive) shocks. On the other hand, with comparing Figure 3(a) and 3(c), it is seen that the intensity of compressive shocks is significant when the plasma model satisfies higher density of hot electron. These results are important and don't see in.³⁷

It is clear that if dissipative effects (due to the kinematic viscosity) in a plasma medium become very strong the shock waves will appear

instead solitons. Therefore, to complete discussion, the effect of kinematic viscosity is shown in Figure 4 for different density cases (a) $\beta = 0.6$, (b) $\beta = 0.85$, and (c) $\beta = 0.995$, respectively. It is observed that ion kinematic viscosity increases the amplitude both compressive and rarefactive shock waves. On the other hand, it is

shown that shock strength increases as the viscosity in the system is increased. It is also clear that transformation from solitons profile to shock structure occur in large values η_0 . The results are similar to that obtained in³⁷ for an oblique propagation of shock wave in a superthermal magnetized plasma.

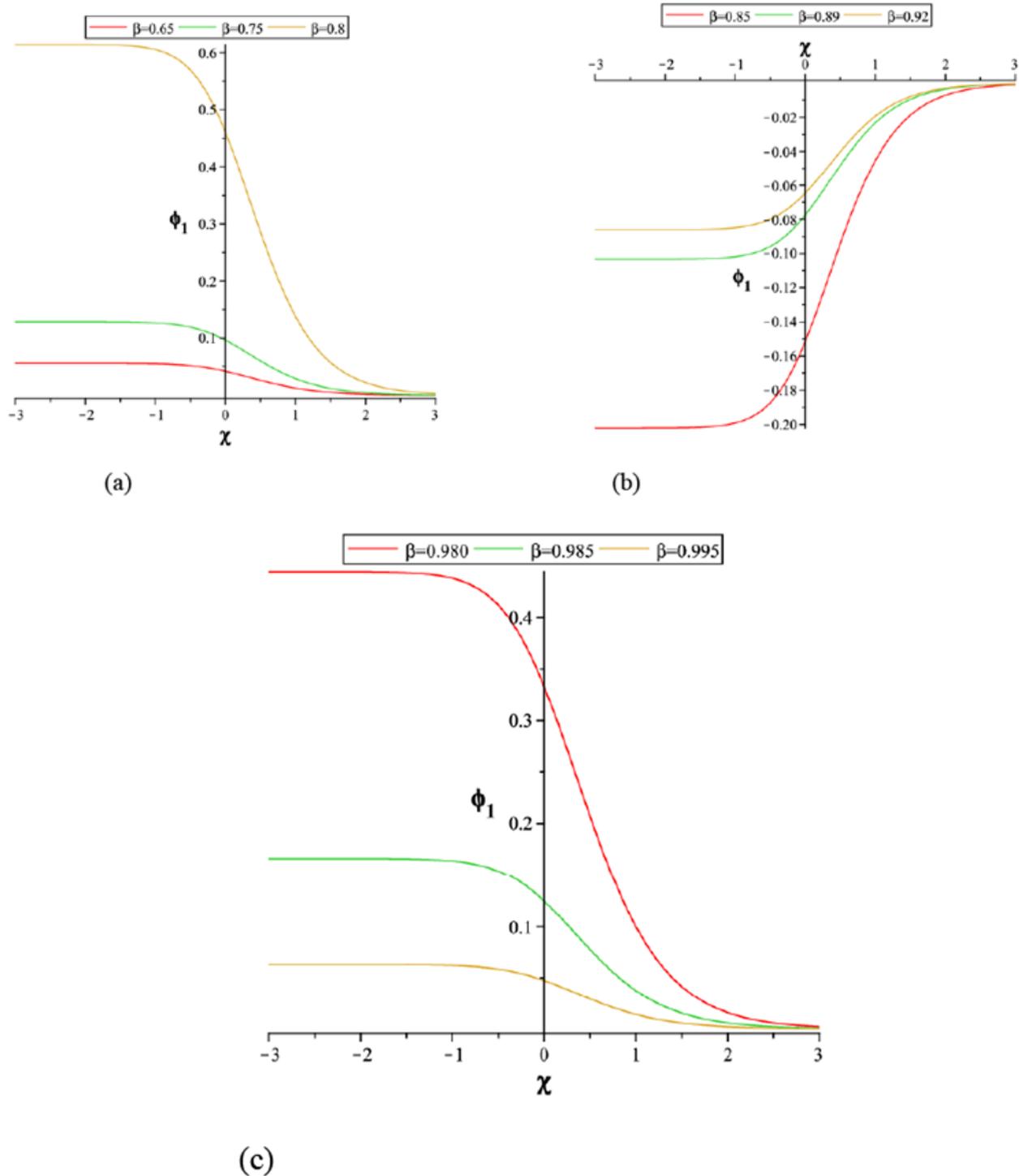


Figure 2 Variation of φ_1 with χ - Eq. (16) - for different values of β with $k=3$, $\tau =20$, and $\eta_0 =0.3$ for (a) $0 < \beta < \beta_{c1}$ ($=0.812$), (b) $\beta_{c1} < \beta < \beta_{c2}$ and (c) $\beta > \beta_{c2}$ ($=0.976$).

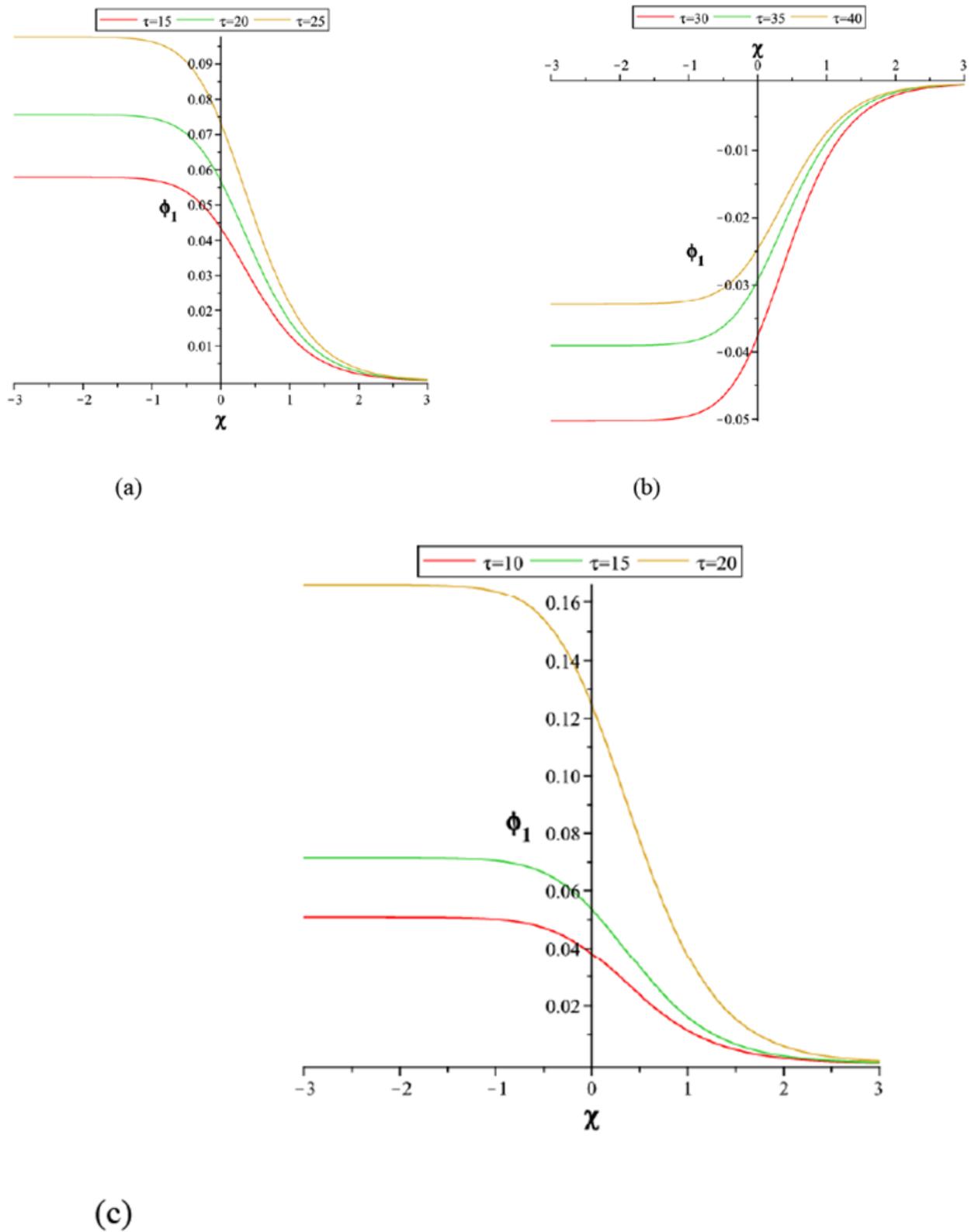


Figure 3 Variation of ϕ_1 with χ - Eq. (16) - for different values τ with $k = 3$ and $\eta_0 = 0.3$ for (a) $\beta = 0.7$ (b) $\beta = 0.85$ and (c) $\beta = 0.985$.

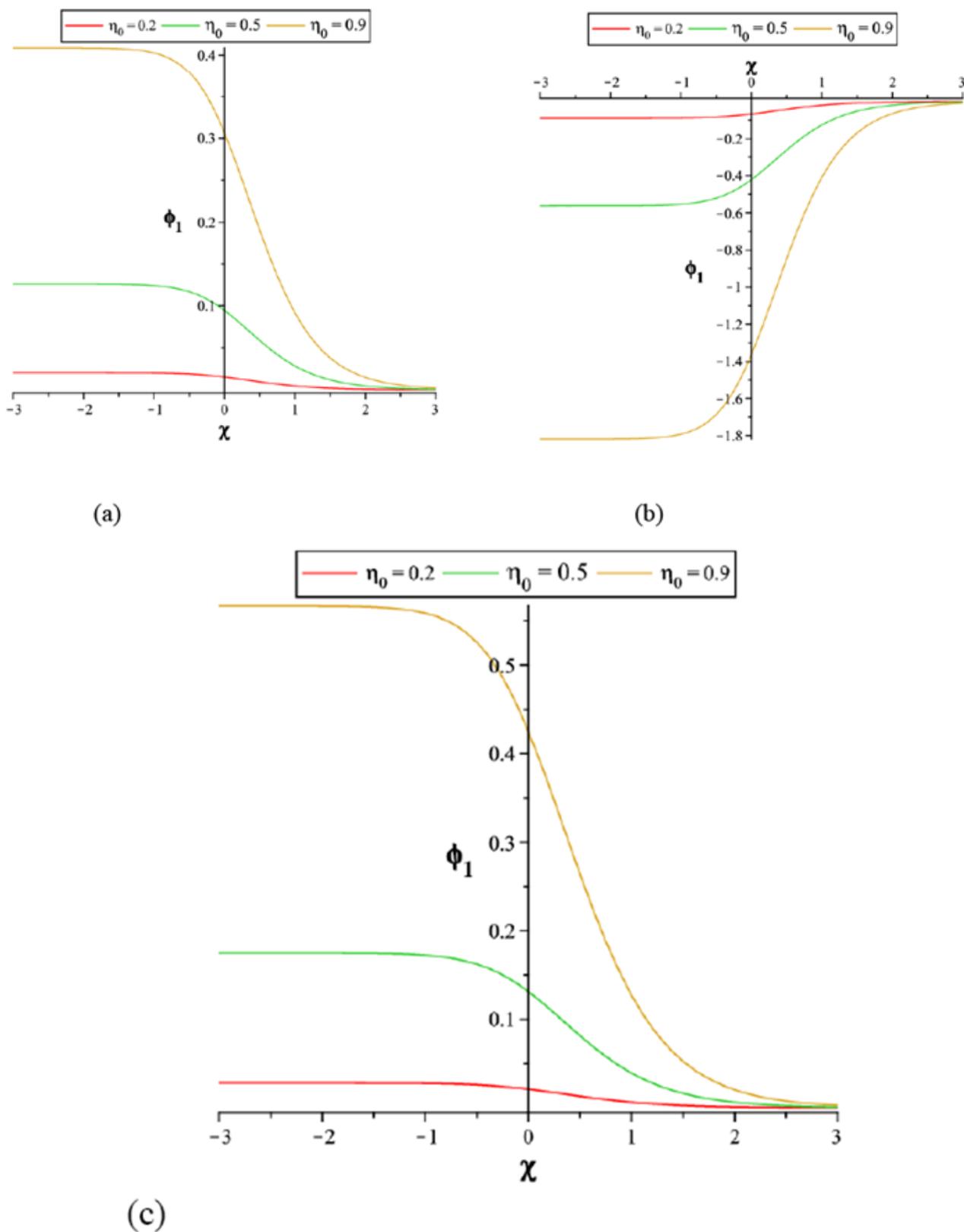


Figure 4 Variation of ϕ_1 with χ - Eq. (16) - for different values of η_0 with $k = 3$ and $\tau = 20$ for (a) $\beta = 0.6$, (b) $\beta = 0.85$ and (c) $\beta = 0.995$.

In continue characteristics of the Burgers shock waves -described by Eq. (17)– is investigated. In Table 1, we give values of the amplitude φm and width Δ shock waves for various values of density β . It is seen that the amplitude of the compressive shock waves increases (decreases) as the density is lower (larger) than β_{c1} (β_{c2}), while the amplitude of the rarefactive shocks decreases by increasing density. In addition, it is interesting to note that the width of the shock waves seems to decreases with density β for all cases of density.

The influence of hot electron temperature τ on the amplitude φm and width Δ shockwaves – described by Burgers equation (17) – is examined in Table 2. This table show that as the hot to cold electron temperature ratio increases the amplitude of compressive shock waves increases. However, for the rare active shock, the amplitude of a shock decreases with increase in τ . Also, it is clearly depicted that the considered plasma model support wider shocks as hot electron temperature increases. We have also investigated the effect of the kinematic viscosity on shock waves propagation in such plasma in Table 3. It is found that an increase in η_0 (meaning an increase in dissipation) considerably caused an increase in the amplitude both compressive and rarefactive shocks. On the other hand, it is found that shock strength increases as the viscosity in the system is increased. In addition, we see from Table 3, that shock waves have low width and propagated narrower for large kinematic viscosity coefficient.

Table 1 The variations of the amplitude φm and width Δ shock waves – described by Burger equation- for fixed values $\tau = T_h / T_c = 20, k = 3\eta_0 = 0.3$ and the various values of β

The different cases for density	β	φm	Δ
$0 < \beta < \beta_{c1}$	0.65	0.023	30.68
	0.75	0.053	29.41
	0.80	0.256	28.46
$\beta_{c1} < \beta < \beta_{c2}$	0.85	-0.084	27.12
	0.90	-0.039	25.14
	0.93	-0.035	23.39
$\beta > \beta_{c2}$	0.98	0.185	18.63
	0.985	0.069	17.94
	0.995	0.026	16.38

Table 2 The variations of the amplitude φm and width Δ shock waves – described by Burgers equation- for fixed values $k = 3\eta_0 = 0.3$ and the different values of τ

The different cases for density	τ	φm	Δ
$\beta = 0.7$	15	0.024	29.3
	20	0.031	30.13
	25	0.04	30.67
$\beta = 0.85$	30	-0.02	28.66
	35	-0.016	29.17
	40	-0.013	29.58
$\beta = 0.985$	10	0.021	16.8
	15	0.029	17.39
	20	0.069	17.94

Table 3 The variations of the amplitude φm and width Δ shock waves – described by Burgers equation- for fixed values $\tau = 20, k = 3$ and the different values of η_0

The different cases for density	η_0	φm	Δ
$\beta = 0.6$	0.2	0.008	46.7
	0.5	0.052	18.68
	0.9	0.17	10.37
$\beta = 0.85$	0.2	-0.037	40.69
	0.5	-0.233	16.27
	0.9	-0.757	9.04
$\beta = 0.995$	0.2	0.011	24.57
	0.5	0.072	9.82
	0.9	0.236	5.46

Conclusion

We have investigated the nonlinear properties of shock waves in a nonlinear and dissipative media consisting of cold fluid ions and cold and hot (superthermal) electrons. Using the reductive perturbation method, the basic set of equations is reduced to K-dVB equation for lowest order perturbation. By assuming quasi-neutrality, Burgers equation is also derived. Exact stationary shock solution of the K-dVB and Burgers equations is presented through the tanh method [42-43] and the results obtained from present model can be summarized as follows:

The nonlinear term A is zero for some critical value of $\beta = \beta_c$, therefore shock waves with positive or negative amplitude can be found in present media.

Increase in hot electron to ion unperturbed number density ratio leads to the increase or decrease in the amplitude of shock waves. On the other hand, when the number density is lower than β_{c1} , the amplitude of shock waves increases, while it decreases for the number density larger than β_{c1} in both cases positive and negative shock waves.

Increase in the ratio of the hot to cold electron temperature τ have also the significant effects on the amplitude of the shock waves.

An increase in the dissipation factor of the homogeneous plasma medium η_0 leads to an increase the amplitude of shock waves.

Increase in the number density of hot electrons β and the kinematic viscosity η_0 lead to decrease in the width of shock waves while increase in the temperature of hot electron τ leads to increase in the width of shock waves.

The findings of this investigation may be important in the laboratory and astrophysical/space plasma environments where two populations of electrons (cold and hot) and cold ions are present. In end, it should be mentioned that the effects warm ions, magnetic field and also the treatment of shock waves exactly in $\beta = \beta_c$ can be considered in such plasmas in future works.

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Conflicts of interest

The author declares there is no conflict of interest.

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