

Research Article





# Turbulence modeling of drag reducing fluid flow using modified k-w model

#### **Abstract**

In the present study, a standard  $k - \omega$  eddy viscosity turbulence model for Newtonian fluids has been developed to predict the drag reduction and the mean velocity profile of drag reducing fluids. In addition, this model is an extension of FENE-P  $k-\omega$  eddy viscosity model for viscoelastic fluids without using any additional constitutive equation. In this work, the first section is concerned with developing a new correlation for the eddy viscosity equation, which depends on the flow and rheological parameters, based on the Direct Numerical Simulation (DNS) data. For this purpose, a User Defined Function (UDF) was interpreted to Fluent 16.1 solver to supersede the Newtonian eddy viscosity with a new average eddy viscosity which is pertinent to the characteristics of the polymer solution. In the second section, the validity of the created model was compared with the results of experiments and DNS model for fully developed turbulent channel flow for other

The developed model predictions for the drag reduction and the mean velocity profile showed good agreement with the DNS model and experimental results at low and intermediate drag reduction regions (i.e.  $DR \le 60$ ). In contrast, in high drag reduction and near Virk's asymptote region (i.e. DR > 60), the accuracy of the model was lost. Finally, the present model is an extension of Reynolds averaged Navier Stokes equation based on standard  $k-\omega$  model. Consequently, the model can be more practical at higher values of Reynolds number and complex geometries other than the DNS model, which poses benefits only at lower values of Reynolds number and simple geometry.

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## Introduction

The addition of a few weight parts per million (wppm) of longchain flexible polymer molecules to the turbulent flow of a solvent can drastically improve the heat transfer and reduce the turbulent friction. This process is also known as a drag reduction. Always, higher molecular weight polymers give better drag reduction performance. However, the addition of a minute amount of polymer in the working fluid reduces the turbulent eddies due to suppressing the formation and propagation turbulent bursts in the buffer layer region. Consequently, resulting in losses reduction in hydraulic energy due to a chaotic and random motion.<sup>1</sup> This phenomenon was initially discovered by.<sup>2</sup> Since then drag reduction by polymer additives subsequently undergoes extensive researches due to continuous increase in its practical applications, such as, long distance fluid transportation and in cooling and heating systems. Therefore, developing the numerical models becomes a vital issue to solve these types of practical engineering problems. To understand the polymer behavior and why drag reduction occur, many experimental and numerical studies had been carried out in the fields of chemistry and fluid mechanics. The exact mechanism of drag reduction for polymer solutions is still vague due to the complexity of flow.

There are two phenomena in drag reducing fluid flow as following:

- Onset of drag reduction.
- Maximum Drag Reduction (MDR).

The onset of drag reduction is the criteria which must be met before a given polymer solution can display any drag reducing effects in turbulent flow. The onset of drag reduction relies on polymer solution parameters and Reynolds number. If these criteria are not met, the flow behaves similar to a Newtonian fluid and the drag reduction is not occurred. The mean velocity profile of Newtonian fluids can be expressed by:

$$u^{+} = y^{+}$$
  $y^{+} < 11.6$  (sub layer region) (1)  
 $u^{+} = 2.5 \ln y^{+} + 5.5$   $y^{+} \ge 11.6$  (buffer layer region) (2)

Where  $u^{+}$  is the normalized velocity i.e.,  $u/u_{\tau}$ , u is the local velocity and  $u_{\tau}$  is the frictional velocity,  $y^{\dagger}$  is the normalized distance from a wall i.e.,  $y^+ = u_{\tau} y / v$ , v is the kinematic viscosity and y is the normal distance from a wall. When the drag reduction value increases, observed that, the log law is shifted by  $S^+$  in buffer layer region and the slope of line is not changed. It was concluded that, the mean velocity profile in buffer layer can be written in the form of:

$$u^{+} = 2.5 \ln y^{+} + 5.5 + S^{+}$$
 (3)

In addition,3 found a universal asymptote which is not related to both of the type and the polymer concentration in the solution. This phenomenon is known as maximum drag reduction (MDR). Based on experimental data,3 proposed the mean velocity profile in buffer layer region at (MDR) by:

$$u^{+} = 11.7 \ln y^{+} - 17 \tag{4}$$

There are three regions of drag reduction as the following:

Low region (DR < 30 %)

Intermediate region (30 %  $\leq DR \leq$  50 %)

Maximum region (DR > 50 %)

Interest in developing turbulence models for the prediction of drag reducing fluids flow has grown over recent years. The main object of drag reducing fluids researches is to develop a viscoelastic turbulence





model which is able to predict correctly the drag reduction value and the mean velocity profile in all drag reduction regions and for any polymer solution. The first turbulence models for viscoelastic fluids flow originated at 1970's by<sup>4</sup> and.<sup>5</sup> The limited range illustrates how complexity, difficulty and weakly of such models, since they depend on large number of parameters that have different values for each polymer solution in each flow situation. Also these earlier models had no obvious link to the polymer rheology. More physicallybased models for purely viscous fluids of variable viscosity were investigated by. 6 Based on the same concept, Pinho and co-workers, 7-9 later adopted a generalized Newtonian fluid constitutive equation by introducing some extensional viscosity influences that include the deformation tensor rate. These models tried to describe the polymer rheology by purely viscous constitutive equation such as the power law model for yield stress fluids. However, these models are not based on a true viscoelastic constitutive and it was unable to capture all features of elastic fluids.

It is well known that, the DNS constitutive models, which describe the dilute polymer solutions rheology well like FENE-P model, were adopted for the turbulence closures development. The FENE-P is the more accurate model for these complex flow situations, regardless that there are still some discrepancies between the measured intensities of drag reduction and calculated (by DNS) specially at higher drag reduction region.<sup>10</sup> Therefore, <sup>11-13</sup>used extensively the FENE-P model, which is coupled with DNS model, to simulate the rheology of the fluids in turbulent flows. However, little number of publications based on Giesekus model<sup>14</sup> and Oldroyd-B model,<sup>15</sup> which are suitable for the surfactant flow, were coupled with DNS model to simulate the FENE-P fluids. In some cases, the DNS data was processed to provide Reynolds average data<sup>16,17</sup> and explain the vortex dynamics.<sup>18,19</sup> Though, there is an absence of detailed experimental evaluation, DNS simulation can provide a validating guidelines for both Reynolds Averaged Navier-Stokes (RANS) equations and Large eddy simulations (LES) at relatively low Reynolds number and simple geometry such as a flat plate flow or a fully developed channel flow. On contrary, at high Reynolds number and complex geometry such as flow past ships, the DNS simulation are impractical, because the number of degree of freedom for DNS  $\propto Re^{9/4}$ 

Consequently, the need to develop turbulence models based on RANS equations by using FENE-P model has been appeared in recent years. 13 developed a Reynolds average kind of turbulence model based on FENE-P models. The mathematical expressions for eddy viscosity of RANS mixing length model in three regions of drag reduction was presented based on the DNS results of the FENE-P model. Then <sup>20</sup> proposed a model that is considered as an extension of the low Reynolds number  $k - \varepsilon$  model of Nagano Y et al.<sup>21</sup> for Newtonian fluids.<sup>20</sup> used a two additional constitutive equation based on FENE-P model and developed closures of terms in conformation tensor equation based on the results of Housiadas KD et al.<sup>18</sup> These earlier models were typically correlated and developed based on the DNS data and were only used to understand and predict the low and intermediate drag reduction regions. 22 developed a  $k - \varepsilon$  model based on the  $v^2 - f$  model which is an extension of Durbin P<sup>23</sup> model. The Reynolds averaged polymer stress was modeled without solving the evolution equation for the polymer conformation tensor. The results showed a good agreement with DNS data while model is still limited to low and intermediate drag reduction regions. More recently,<sup>24</sup> improved significantly the turbulence model of Pinho FT et al<sup>20</sup> by extending its closures of terms in conformation tensor equation to the high drag reduction region and improving its general performance. For Newtonian fluids, it is known that there are problems in the  $k-\varepsilon$  turbulence models either because of the lack of natural boundary condition of,  $\varepsilon$ , and the appearance of higher-order correlations in the balance of the dissipation rate at the wall. As a result, forcing the usage of higher-order derivatives of the turbulent kinetic energy that leads to a not asymptotic behavior and consequently to numerical stiffness. So researchers tried to find an alternative to  $\varepsilon$ . Wilcox developed the  $k-\omega$  model where the exact viscous terms next to the wall do not require modeling. Thus leading to better predictions without using the damping functions typical of  $k-\varepsilon$  closures.

Consequently,<sup>25</sup> developed an eddy viscosity model based on<sup>26</sup>  $k-\omega$  model for drag reducing fluid flow. The new eddy viscosity model based on polymer rheology is an extension of Pinho FT et al<sup>20</sup> and two additional constitutive equation were adopted. These models are typically used for drag reduction up to 50 % in low and intermediate region. Another approach to turbulence modeling was proposed by<sup>27</sup> which is an extension of Iaccarino G et al<sup>22</sup> but it is valid for the three regions of the drag reduction. The additional constitutive equation based on FENE-P model is not used in this model and it's replaced with two transport equations for  $v^2$  and f. This model was modified by<sup>28</sup> to predict the mean temperature profile and heat transfer rates in viscoelastic fluid by adding the additional terms to the energy equation. Moreover,<sup>29</sup> modified the model developed by<sup>27</sup>  $(k - \varepsilon - v^2 - f \mod el)$  to predict the velocity profiles, shear stress profiles and the percentage of friction drag reduction at high concentration solutions. The accuracy of model was lost at both of high Reynolds number and high concentration.

Based on the previous review of literature, it has been concluded that:

- a) The DNS constitutive models (like FENE-P model), which describe well the polymer rheology, are more accurate model for the turbulent drag reducing fluid flow at low values of Reynolds number and simple geometries.
- b) At higher Reynolds number values and complex geometries, the RANS models are more practical than the DNS model.

For wall bounded turbulent flow, better predictions were given by the RANS  $k-\omega$  turbulence model without using additional damping function as in the case of using  $k-\varepsilon$  model.

Using additional FENE-P constitutive equations in RANS  $k-\omega$  model increase the running time and simulation cost, especially at higher values of Reynolds number and complex geometries.

Hence, the objective of the present study is to investigate a new RANS  $k-\omega$  turbulence model for drag reducing flow without using additional constitutive equations. The investigated model is based on the standard RANS  $k-\omega$  turbulence model developed by <sup>26</sup> and eddy viscosity model proposed by <sup>25</sup> A user defined function (UDF) code is interpreted to ANSYS -Fluent 16.1 solver to replace the Newtonian eddy viscosity with a new average eddy viscosity, which takes into account the flow characteristics and the polymer parameters. The predicted model will be valid for both low and intermediate drag reduction regions ( $DR \le 60$ ), because of the limiting assumption of Resende model. The model also will be verified using DNS data for fully-developed channel flow provided by Li C et al<sup>13</sup> and the experimental Results of Warholic M et al.<sup>30</sup>

## **Numerical** model

#### **Governing equations**

The turbulence model developed here is based on the continuity and the Reynolds Averaged Navier Stokes (RANS) equations.

## Incompressible continuity and momentum equations

The equations are written in the indicial notation of Einstein with Kronecker de  $(\delta_{ij} = 1 \text{ if } i = j \text{ and } \delta_{ij} = 0 \text{ if } i \neq j)$ . The time-averaged continuity and momentum equations for an incompressible FENE-P fluid are:

$$\frac{-i}{\partial x} = 0$$

$$\frac{\partial U}{\partial t} = 0$$

$$\frac{\partial U}{\partial t$$

Where  $U_i$  is the mean averaged velocity, P is mean averaged pressure,  $\rho$  is the fluid density,  $\overline{\tau}_{ik,p}$  is mean averaged polymer stresses,  $-\rho \overline{u_i u_k}$  is the Reynolds stress tensor and  $\mu$  is the viscosity of the Newtonian solvent. The capital letters and overbars denote Reynolds-averaged quantities, whereas small letters denote fluctuations.

## Wilcox $k - \omega$ turbulence model

The eddy viscosity (  $\mu_{t}$  ) for Newtonian fluid is given by :

$$\mu_t = \rho \frac{k}{\omega} \tag{7}$$

The Reynolds stresses are computed as usual in two-equation models with the Boussinesq expression in the form of:

$$-\rho \overline{u_i u_k} = 2\rho v_T S_{ij} - \frac{2}{3}\rho k \delta_{ij}$$
 (8)

where k is turbulence kinetic energy,  $\omega$  is the turbulence frequency.

The transport equations for k and  $\omega$  for turbulent flow can be expressed by Eqs (9&10)

$$\frac{\partial \rho k}{\partial t} + \frac{\partial \rho U k}{\partial x} = \frac{\partial}{\partial x} \left[ \begin{pmatrix} \mu \\ \mu + \frac{t}{s} \\ s \\ \kappa \end{pmatrix} \frac{\partial k}{\partial x} \right] + \begin{pmatrix} \partial U \\ 2\mu S . S - \frac{2}{3}\rho k \frac{i}{\delta x} \\ s \\ i \\ j \end{pmatrix} - \beta * \rho k \omega$$
(9)

$$\frac{\partial \rho \omega}{\partial t} + \frac{\partial \rho U \omega}{\partial x} = \frac{\partial}{\partial x} \left[ \begin{pmatrix} \mu \\ \mu + \frac{i}{s} \\ s \\ \sigma \\ \omega \end{pmatrix} \frac{\partial \omega}{\partial x} \right] + \gamma_{1} \begin{pmatrix} 2\rho S \cdot S - \frac{2}{3}\rho \omega \frac{i}{\partial x} \\ i & j \end{pmatrix} - \beta_{1}\rho \omega^{2}$$

$$(10)$$

The model constants are  $\sigma_k = 2.0$ ,  $\beta^* = 0.09$ ,  $\sigma_m = 2.0$ ,

$$\gamma_1 = 0.553$$
,  $\beta_1 = 0.075$ .

#### Non-dimensional parameters

In present work, the following non-dimensional numbers are used, which are defined as:

Frictional Reynolds number  $Re^{\tau} = hu^{\tau}/v^{o}$ , where  $h^{t}$  is the channel half height and  $v^{t}$  is the kinematic viscosity of the polymer and solvent ( $v_{o} = v_{p} + v_{s}^{o}$ ).

**Weissenberg number**  $(We_{\tau o} = \lambda u_{\tau}^2 / v_o)$  is the ratio of the polymer relaxation time,  $\lambda$ , to the shortest time-scale,  $v_o / u_{\tau}^2$  in the viscoelastic turbulent flow.

Viscosity turbulent ratio  $\beta = v_s / v_o$  is the ratio of the solvent viscosity to the sum of the kinematic viscosities of the solvent and polymer.

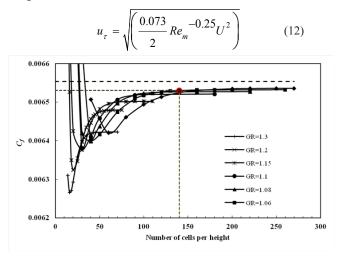
**Drag reduction ratio** (DR) is the ratio of the reduction of the wall shear stress in viscoelastic flow  $(\tau_{w,v})$  to the Newtonian wall shear

$$DR = \frac{\tau_{w,n} - \tau_{w,v}}{\tau_{w,n}} = 1 - \frac{\tau_{w,v}}{\tau_{w,n}}$$
(11)

#### Mesh independence test and validation

The fluid domain was created using ANSYS 16.1 Design Modeler for channel with half height, h of 5 cm and 3 m length. The mesh independence test was performed and created using ANSYS-Mesh 16.1 with range of (10 to 270) cells non uniformly distributed across the channel height with growth rate (GR) from (1.3 to 1.04). The simulation was performed for frictional Reynolds number,  $Re_{\tau} = 395$  where this value corresponds to mean flow Reynolds number.  $Re_m = 13850$ . The mean flow Reynolds number is based on the channel height,  $Re_m = 2hU_m / v_o$ , where  $U_m$  is the mean flow velocity. Figure 1 shows the change in the skin coefficient  $C_{\ell}$  at the end of channel for different mesh sizes. The figure also illustrates that, the skin coefficient decreases with increasing the number of cells until it reaches to the inflection point and hence increasing to asymptotic value. Figure 2 reveals the change in skin coefficient  $(C_d)$  with the value of first cell  $y^{+}$ . It is depicted that, the skin coefficient decreases with decreasing the first cell height and hence increases until it reaches asymptotic value. It was clearly seen that with decreasing the value of growth rate, the number of cells to yield constant value of  $C_{\ell}$ increases, and hence the first cell height  $y^{+}$  is decreased (i.e. to reach a constant value of  $C_\rho$  the number of cells per height increases with decreasing the growth rate).

Nevertheless, the relation between the relative error and the number of cells required to reach a constant value of  $C_{\epsilon}$  at each growth rate is represented in Figure 3 to obtain the optimum growth rate for the present study. The figure shows that the growth rate less than 1.06 with relative error less than 0.36 % reaches to approximately a constant value of relative error. Therefore, from the previous study, the optimum growth rate and value of first cell  $y^{+}$  at the wall were chosen as 1.06 and 0.00378, respectively. Furthermore, the frictional velocity and hence the skin coefficient were compared with experimental data from Dean's correlation according to the literature <sup>31</sup> as given by Eq.(12). Also, the final results of the present model are compared with the universal log law for velocity profile as illustrated in Figure 4.



**Figure 1** The change of skin coefficient with number of cells per height of the channel at different growth rate values.

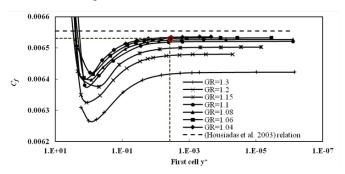


Figure 2 The change of skin coefficient with the first cell y+ at different growth rate values.

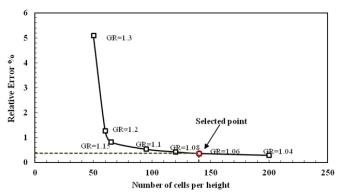


Figure 3 The change of relative error with number of cells to yield constant value of  $\,C_f\,$  according to each growth rate.

It is observed from Figure 4 that, the results for the present model agree well with the experimental data which is taken from Dean's correlation method and both the two curves are indistinguishable on the scale of the figure except for moderate values of  $y^+$ . The error at number of cells per height equals 140 is sufficiently small so that this value is considered acceptable for the present study.

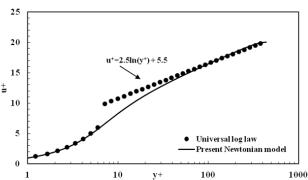


Figure 4 Velocity profile for turbulent channel flow compared with experimental correlation.

# Development of approximated $k - \omega$ model

# Traditional model methodology

For turbulent flow, it is more usually to model the Reynolds stresses by Boussinesq approach as described in Eq.(8).  $^{26}$  computed the Reynolds stresses by two-equation models. He introduced the transport equation for k and  $\omega$  as represented by Eqs.(9) and (10) for Newtonian fluid. Despite of eddy viscosity for Newtonian fluid is expressed by Eq.(7), it is the most challenge to address for viscoelastic fluids turbulence modeling. Therefore, for drag reducing fluid, the selection of the most adequate expression is of relevant importance to improve the turbulence detection. The most commonly used method for identifying the kinematic eddy viscosity,  $\nu_T$  was developed by  $^{25}$ . They modified the eddy viscosity model to be capable of predicting the drag reduction of viscoelastic fluids at different values of Weissenberg mumbers. In addition, Resende adapted the kinematic eddy viscosity,  $\nu_T$  to be the sum of a Newtonian and a polymeric contribution as given by Eq.(13).

$$v_T = v_T^N - v_T^p \tag{13}$$

The polymeric contribution to the eddy viscosity,  $v_T^p$  is dependent on the trace of the conformation tensor,  $C_{mm}$  and introduces a damping function which is relevance to both of Weissenberg number and distance from the wall, i.e.,  $f\left(We_{r_a}, y^+\right)$ .

$$v_T^p = f\left(We_{\tau_0}, y^+\right) \times \frac{C_{mm}}{\sqrt{L^2}} \times C_\mu \times f_\mu \times \frac{k}{\omega^N}$$
 (14)

Where:

$$f\left(We_{\tau_{0}}, y^{+}\right) = C_{u}^{p} \times f_{u}^{p} \times f_{DR}$$
 (15)

$$C_{\mu}^{p} = 0.0135 \times \left[ \frac{25}{We_{\perp}} \right]^{0.12}$$
 (16)

$$f_{\mu}^{p} = \left[ 1 + 2.55 \times \exp\left(-\frac{y^{+}}{44}\right) \right]$$
 (17)

$$f_{DR} = \left[ 1 - \exp\left( -\frac{We_{\tau o}}{6.25} \right) \right]^4 \tag{18}$$

# New methodology description

In the present work, a new average eddy viscosity is supposed,  $\tilde{\mathcal{V}}_T$  (based on Resende model), is derived and written in the following

expression:

$$\tilde{v}_T = v_T^N \left( 1 - \frac{v_T^p}{v_T^N} \right) \tag{19}$$

<sup>17</sup>approved that there are four parameters that influence DR namely, the Weissenberg number  $(We_{\tau_o})$ , the molecular flexibility  $(L^2)$ , the viscosity Ratio  $(\beta)$  and the frictional Reynolds number  $(Re_{\tau})$ . Consequently, the trace of conformation tensor in Resende damping function is replaced with a new function  $(M_{we})$  which has a direct dependency on the previous four parameters. Therefore, the new derived average eddy viscosity  $\tilde{v}_T$  can be expressed as the form of Eq.(14) and is presented by Eq. (20). The new supposed average eddy viscosity  $(\tilde{v}_T)$  is set to depend on both of Resende function,  $f(We_{\tau_o}, y^+)$ , and  $M_{we}$  function.

$$\tilde{v}_T = v_T^N \left( 1 - f\left( W e_{\tau o}, y^+ \right) \times M_{we} \right) \quad (20)$$

# Enhanced new $M_{we}$ function and UDF code verification

The values of  $M_{we}$  are calculated at different values of maximum extensibility parameter (L) at three different values equal to 25, 50, and 100 of Weissenberg number  $(We_{\tau_0})$ . Moreover, the results are explored as well to three different  $Re_{\tau}$ , values equal to 125, 180, and 395 for dilute solution  $(\beta \simeq 1)$ . The values of  $M_{we}$  are selected to

obtain the same values of DNS drag reduction modeling by using Resende eddy viscosity Model. UDF code is interpreted to ANSYSFluent 16.1 solver and Newtonian turbulent viscosity replaced with average turbulent viscosity as represented by Eq.(20). Figure 5 shows the variations of  $M_{\rm WE}$  values with the maximum extensibility parameter, L, for three various values of  $W_{\rm E_{70}}$  equal to 25, 50, and 150 at three different values of  $W_{\rm E_{70}}$  and  $R_{\rm E_{7}}$  the value of  $M_{\rm we}$  increases with increasing L and eventually asymptotes. Figure 5 also shows that, the effect of varying the value of  $M_{\rm we}$  increases with increasing  $W_{\rm E_{70}}$ , and it decreases with increasing of  $R_{\rm E_{70}}$  at the same value of L and  $W_{\rm E_{70}}$ . Consequently, Figure 6 shows the fitting data, in dashed line, by the Least Square Method with minimum value of determination factor,  $R^2 \simeq 0.93$  over the entire range of  $M_{\rm we}$  values resulting in Eq.(21) with the relevant constants values as following:

$$M_{we} = a \left( 1 - \exp\left(-bL\right) \right) \times \left( 1 - \exp\left(-c\left(We_{\tau o} - d\right) \left(\frac{Re_{\tau}}{125}\right)^{-e}\right) \right)$$
(21)

Where:

$$a = 386.336$$
;  $b = 0.0425$ ;  $c = 0.0247$ ;  $d = 0.0986$ ;  $e = 0.1157$ 

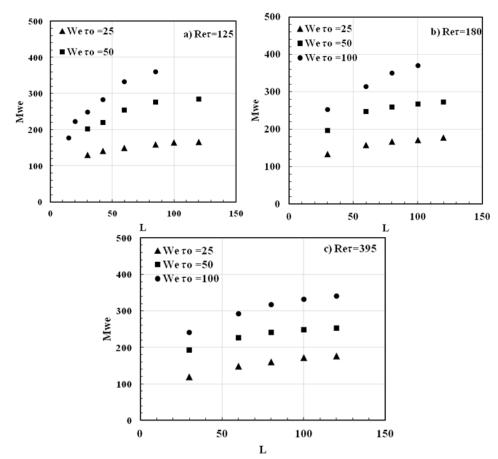


Figure 5 Variation of  $M_{we}$  with maximum extensibility parameter, L, and Weissen berg number equal to 25, 50, 100 at a:  $Re_{\tau}=125$ ; b:  $Re_{\tau}=180$ ; c s:  $Re_{\tau}=395$ .

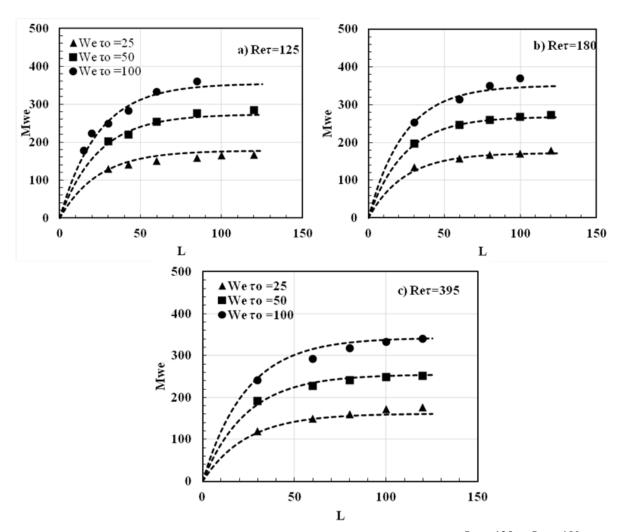


Figure 6 Fitting data of Mwe function with maximum extensibility parameter, L, and Weissenberg number at a:  $Re_{\tau}=125$  ;b:  $Re_{\tau}=180$  ;c:  $Re_{\tau}=395$  .

### **Results and discussion**

The new proposed turbulence model was validated by DNS model for viscoelastic FENE-P fluids in fully developed turbulent channel flow. The simulations were performed for different values of frictional Reynolds number  $(Re_{\tau})$  maximum extensibility parameter of polymer (L) and different values of Weissenberg number  $(We_{\tau_0})$ . The new investigated model predictions for both of mean velocity profile and turbulent viscosity ratio were compared with the DNS model, Resende  $k - \omega$  turbulence model and other RANS eddy viscosity models at different values of frictional Reynolds number, Weissenberg number and maximum extension of polymers. Furthermore, the model compared with the experimental results of Warholic M et al.<sup>30</sup> at different values of polymer concentration in low, intermediate and higher drag reduction regions for fully-developed turbulent channel

# The developed model results validation with DNS

The results of the present turbulence model were validated with the DNS results of Li C et al13 which are based on FENE-P model at all regions of drag reduction (i.e. low, intermediate and high drag reduction). As seen in Table 1, the present model predictions show

excellent agreement with the DNS results at low and intermediate drag reduction regions (i.e.  $DR \le 60 \%$ ) with a maximum error only reaches about 5 %. However, the error increased in some cases due to the discrepancy between the fitted equation points and the actual value of drag reduction such as the case of  $Re_{\tau} = 395$ ,  $We_{\tau 0} = 100$  and L=60. The discrepancy appears clearly in Figure 7. On the other hand, in the case of high drag reduction region and near Virk's asymptote, the accuracy of the model is lost as expected.

Typical mean velocity profiles as a function of the normalized distance from a wall  $(y^+)$  in logarithmic scale in the low, intermediate and high drag reduction regions are shown in Figures 7-9. For the sake of comparison, the mean velocity profiles for Newtonian flow have been included at every Reynolds number.

For the Newtonian cases, excellent agreement was obtained for the velocity distribution in the viscous sublayer  $(u^+ = y^+)$  as well as in the log-law layer  $(u^+ = 2.5 \ln(y^+) + 5.5)$ . It can be seen that, all the velocity profiles in the drag reduced flow have the same linear distribution  $(u^+ = y^+)$  in the viscous sub layer.

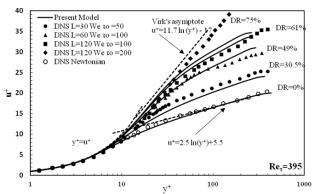


Figure 7 Mean velocity profiles predictions compared with the DNS data of 13 for turbulent channel flow with  $Re_\tau=395$  .

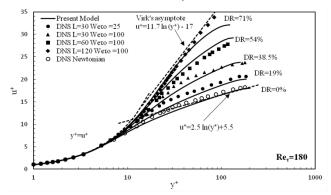


Figure 8 Mean velocity profiles predictions compared with the DNS data of 13 for turbulent channel flow with  $Re_{\tau}=180$  .

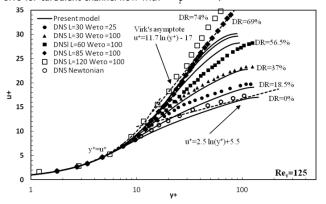


Figure 9 Mean velocity profiles predictions compared with the DNS data of 13 for turbulent channel flow with  $Re_z=125$  .

Further away from the wall, specifically in the logarithmic layer, the mean velocity profile increases as compared to that in Newtonian flow. In the low drag reduction region, the profile is shifted upwards with the same slope of Newtonian profile, which is consistent with the experimental results of Warholic M et al.<sup>30</sup> In the intermediate and higher drag reduction regions, the slope of the velocity profile in logarithmic layer is increased with increasing the value of the drag reduction until it reaches to the Virk's asymptote line.

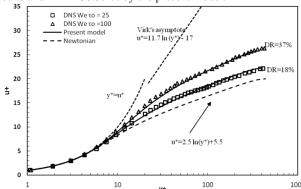
### Results of the model against DNS and resende model

The present model was compared with the DNS of Li C et all<sup>13</sup> and <sup>25</sup> eddy viscosity models within the turbulent-flow regime at

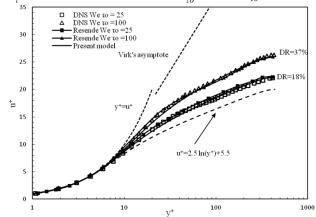
approximately identical frictional Reynolds number  $Re_{\tau}=395$  , L=30 and  $\beta=0.9$  .

# Predicted velocity profile against DNS results and Resende $k-\omega$ model

Figure 10 demonstrates the predicted data of the mean velocity profile against the DNS results of Li CF et al<sup>17</sup> within the turbulentflow regime at approximately identical frictional Reynolds number  $Re_{\tau} = 395$ , L = 30 and  $\beta = 0.9$ . The present model predicted DR = 17.1 % against DR = 18 % by DNS results at  $We_{\tau_0} = 25$ . Also, at  $We_{\tau o} = 100$ , the present model predicted DR = 38.35 %against DR = 37 % by DNS results. The present model predictions show good agreement with the DNS data specially for the trend of the velocity profile and the values of the drag reduction. The monotonic shift of the log-law region with DR is well captured, as well as the correct evolution in the buffer-layer. In Figure 11, the predicted data of mean velocity profile is compared with the results of both of Resende P et al.25 and DNS models. The Resende model predicted DR = 18 % against DR = 18 % by DNS and DR = 17 % by the present model results at  $We_{\tau_0} = 25$ . Also, at  $We_{\tau_0} = 100$  the Resende model predicted DR = 39 % against DR = 37 % by DNS results and DR = 38.35 % by the present model.



**Figure 10** Comparison between the predictions of the mean velocity by the present model and DNS data of 13 for turbulent channel flow with  $Re_{_{7}}=395$ , L=30 and  $\beta=0.9$  at  $We_{_{70}}=25$  &  $We_{_{70}}=100$ .



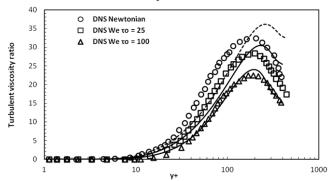
**Figure 11** Comparison between the predictions of the mean velocity by the present model and both DNS model and RANS  $\stackrel{k}{k}-\omega$  model of25 for turbulent channel flow with  $Re_{\tau}=395$ , L=30 and  $\beta=0.9$  at  $We_{\tau\alpha}=25$  and  $We_{\tau\alpha}=100$ .

Table 1 Comparison between the drag reduction predicted by the present model and DNS drag reduction of Li C et al. 13

$Re_{_{ au}}$	L	$We_{ au o}$	DNS DR %	$\mathbf{M}_{\mathrm{we}}$	Model DR %	Relative error %
125	30	25	18.5	128.23	18.23	1.46
	30	100	37	255.1	38.4	3.78
	60	100	56.5	326.3	56.2	0.53
	85	100	69	344.2	59.6	13.62
	120	100	74	351.7	60.7	17.97
180	30	25	19	124.4	18.1	4.74
	30	100	38.5	252.6	38.73	0.6
	60	100	54	323.068	56.67	4.94
	120	100	71	348.23	63.29	10.86
395	30	50	30.5	184.179	29.12	4.52
	60	100	49	315.46	55	12.24
	120	100	61	340.032	61.6	0.98
	120	200	75	378.97	63.7	15.07

# Predicted eddy viscosity ratio against the DNS and resende $k - \omega$ model

As seen from Figure 12, the change of the eddy viscosity ratio  $(v_T/v_o)$  along  $y^+$  is captured especially in the buffer layer region. There is a strong dependency and good agreement of the enhanced predicted eddy viscosities for all studied  $We_{\tau_o}$  values, with DNS eddy viscosities for all  $y^+$  range values. The present model is capable to predict the correct evolution shown by the DNS model in the lower and intermediate drag reduction regimes. Furthermore, a DR increment and turbulent viscosity ratio decrement were achieved.



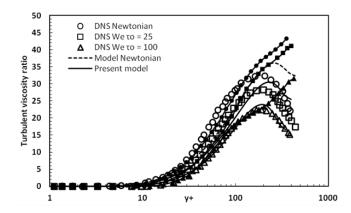
**Figure 12** Comparison between the predictions of the turbulent viscosity ratio by the present model predictions and DNS data of  $^{13}$  for turbulent channel flow with  $Re_{_{\rm T}}=395$  , L=30 and  $\beta=0.9$  at  $We_{_{{\rm T}o}}=25$  and  $We_{_{{\rm T}o}}=100$  .

In Figure 13, the predicted turbulent viscosities were compared with DNS and Resende results. This figure also depicted that there is a marked contrast with Resende model for higher  $y^+$  values for the same Reynolds number. Therefore, it is clear that the present model predictions coincide with those of Resende  $k - \omega$  model and show good agreement with the DNS data.

#### Test the present model against experimental results

The present model was tested against the experimental data of

Warholic M et al<sup>30</sup> at different values of the polymer concentration. The channel used in the experiments was with a height of 5.08 cm and 3 m length. The master polymer was a solution of Percol-727, which is a copolymer of polyacrylamide and sodium acrylamide. The simulations were performed at the same dimensions and frictional Reynolds number  $(Re_{\tau} = 778.4)$ . The polymer parameters for the simulations were calculated by<sup>32</sup> based on<sup>33</sup> experimental results. The mean velocity profile and DR were computed and compared with the experimental results at the three regions of drag reduction (i.e. low, intermediate and high regions). Figure 14 shows a good prediction of the mean velocity profile and the positive shift of the buffer layer with DR, which is well captured in low and intermediate drag reduction regions. At high drag reduction region, the deviation between predicted numerical and experimental mean velocity profile is very large and increased with increasing DR. Table 2 shows the predicted DR variation with the polymer concentration changes in three regions. Also, the relative error with the experimental results is listed in Table 2.



**Figure 13** Comparison between the predictions of the turbulent viscosity ratio by the present model predictions and both DNS model and RANS  $k-\omega$  model of25 for turbulent channel flow with  $Re_{\tau}=395$  , L=30 and  $\beta=0.9$  at  $We_{\tau o}=25$  and  $We_{\tau o}=100$ .

Table 2 Comparison between the drag reduction predicted by the present model and the experimental data of Warholic M et al.30

Concentration (ppm)	$We_{ au o}$	Experimental DR	Model DR	Relative error %
0.24	16.5	14	13.1	6.4
1.25	36.5	33	34.6	4.8
3	74	55	50.5	7.5
50	100	69	59.7	13.5

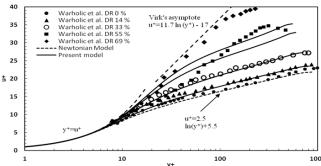


Figure 14 Comparison between the predicted mean velocity by the present model and the experimental data of.<sup>30</sup>

## **Conclusions**

In the present study, the standard  $k - \omega$  eddy viscosity turbulence model has been investigated for the turbulent drag reducing fluid flow. This model has been developed to predict the drag reduction and the mean velocity profile in low and intermediate regions without using any additional constitutive equations. Moreover, an expression for a new average eddy viscosity equation, which is directly related to the fluid and flow rheological parameters, has been developed. In the present model, the trace of the conformation tensor in FENE-P eddy viscosity models was replaced with a new  $M_{we}$  function. This function involves the polymer parameters which affect in the value of the drag reduction. The DNS results were used to predict the  $M_{we}$  function using ANSYS-fluent 16.1 software. A UDF code was interpreted to ANSYS-Fluent 16.1 solver and the Newtonian eddy viscosity was replaced with a new average eddy viscosity equation. The developed model was initially validated with the DNS data and the results of different RANS turbulence models for different values of the frictional Reynolds number and polymer parameters. Also, the developed model was validated by the experimental results at the same flow conditions in low, intermediate and high drag reduction regions. The results obtained from this study can be summarized as follows:

Good agreements between the model predictions and both of the DNS and the experimental results are obtained in low and intermediate drag reduction regions (DR < 60%).

The predicted turbulent viscosity ratio agrees well with the DNS results at low and intermediate drag reduction regions.

The developed model is an extension of the standard RANS  $k-\omega$  model. Consequently, the model could be used at higher values of Reynolds number and complex geometries, which becomes more practical than the DNS model.

The model is simulated without any additional constitutive equations. So, the model reduces the running time and the cost of simulation.

The model also can be used to predict the DR and the mean velocity profile in the developing zone, which may contain heterogeneous polymer solutions.

## **Acknowledgments**

None.

### **Conflicts of interest**

The authors declare that there is no conflict of interest.

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